Monte Carlo Methods for Including Calibration **Uncertainties in Model Fitting Analyses**

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음 0.98 4 Energy [keV] Figure 2. Left: Illustration of the relative

change in the HRMA effectiv

SUMMARY

· Instrument response uncertainties almost universally ignored in astrophysical X-ray data analyses, yet for good quality observations can be dominant source of error

· Response uncertainties are correlated; both understanding and specifying the uncertainties is technically challenging. Moreover, there is no standard set of procedures for incorporating complicated correlated systematic uncertainties in non-linear parameter estimation (eg XSPEC fitting): the approaches used for treating independent errors simply do not apply.

· We have developed Monte Carlo methods to treat calibration uncertainties for the Chandra High Resolution Mirror Assembly (HRMA) and Advanced CCD Imaging Spectrograph (ACIS). The code and ancilliary data will be released to Chandra Users upon acceptance for publication of the article describing this work. CIAO Sherpa methods are also under development to utilise these techniques (see accompanying poster by Kashyap et al)



METHODS

Construct different realisations of instrument response by a combination of (1) randomly varying input parameters describing subassembly performance and (2) random multiplicative perturbation functions, $\mu(E)$, designed to sample subassembly responses with their assessed uncertainties (Fig. 1). Adopt "curtailed Gaussian" probability distribution $P(\sigma)$ for Monte Carlo draws (Fig. 1a). The different subassemblies were treated as follows

HRMA On-Axis: A combination of perturbation functions, $\mu_H(E)$, running within prominent Ir edges, and raytrace-derived effective areas sampling the effects of different hydrocarbon contamination layers with the measured range of allowed values (Fig. 2).

HRMA Vignetting Function: For off-axis angles θ (in arcmin) the area perturbation function was multiplied by a combination of a fractional uncertainty of the azimuthally-averaged vignetting function, $\bar{V}(\theta),$ and an expression involving the ratio of Debye-Waller to perfect mirror reflectivities:

 $\mu_v(E,\theta) = P(\sigma_v)(1-V(\theta)) + \theta P(\sigma_s)(1-R_{DW}/R); \ \sigma_v, \sigma_s = 0.2.$

ACIS OBF and Contamination Layer: OBF transmittance uncertainty employed perturbation functions, $\mu_{OBF}(E)$, between C-N and N-O edges, and O-10 keV constrained by different allowed maximum deviations and relative edge transmittance discontinuity errors. The contamination perturbation function is: $\mu_{CL}(E) = e^{-P(\sigma_C)\tau_C + P(\sigma_C)\tau_O + P(\sigma_F)\tau_F + P(\sigma_F)\tau_{FI}}; \mu_{CL}(0.7keV) < 0.05$

where σ_C , σ_O , σ_F and σ_{Fl} are the fractional uncertainties in the optical depths C, O, F and Fluffium at a fiducial date (2003.29).

ACIS QE: QE uncertainty base on combination of perturbation functions, $\mu_{QE}(E)$, with boundaries at O and Si K edges and ACIS QE model predictions for uncertainties of 13% in CCD depletion depth and 20% in SiO₂ thickness

ACIS Gain and Pulse Height Distribution: Program calcrmf2 used to generate CCD Gain and pulse height response matrix files (RMF) for $P(\sigma_G)$ variations in gain and pulse height width; σ_G =1% @0.7 keV, 0.5% @1.5 keV, and 0.2% @2 4 keV.

ESTIMATING EFECTS OF CALIBRATION UNCERTAINTIES

Perturb nominal effective area, A(E), with combined perturbation functions, $\mu(E)$, to make $A\prime(E)$ (Fig. 3):

$A'(E) = \mu_H(E)\mu_V(E)\mu_{CL}(E)\mu_{OBF}(E)\mu_{QE}(E)A(E)$

For each effective area and RMF, find the parameters of the best-fit model for a synthetic Chandra ACIS observation computed using the nominal instrument response. Compare with parameters found from fits to 1000 synthetic spectra differing only by Poisson noise and generated using the nominal area and RMF (Figs. 4 and 5). This step utilized the XSPEC fitting engine and investigated Blackbody, optically-thin thermal plasma, and power law continuum models. Repeat 1000 times.

Limiting accuracy of *Chandra* ACIS reached in spectra with $\sim 10^4$ counts. Beyond this, errors in in best-fit parameters due to calibration uncertainties completely dominate those due to photon noise.



Figure 4. Example frequency distributions of best-fit parameters obtained for typical blackbody, thermal plasma and powerlaw models from XSPEC for synthetic data sets containing 10⁴ (upper panels) and 10⁵ (lower panels). Black histograms are distributions resulting from 1000 blonce Carlo samplings of the synthetic data allowing Poisson noise writing shows a single and histograms are the distributions examples resulting row fitted single and histograms are the distributions are single and and the single and the single and histograms are the distributions and an anomaly methods from the single and histograms are the ributions of parameters resulting from fits to a single synthetic data set using 1000 Monte Carlo-generated effective areas and response matrice

> Figure 5. Modes and highest posterior density $\pm 96\%$ confidence intervals obtained for the blackbody, thermal plasma and powerlaw models investigated using XSPEC for synthetic data sets containing 10° counts. The y-axes correspond in all cases to the ratio of the input parameter to that retrieved in the prodel fit. Black error bars correspond to 1000 Monte Carlo samplings of the synthetic data and show the effects of Poisson noise variations alone. Dashed error bars correspond to fits to a single synthetic data set using 1000 Monte Carlo-generated effective areas and response matrices.

> > 1.2

0.8

0.6

1.5

0.5

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Absorbed Power Law

▲ n_H=10²

n = 10²

44

<u>.</u>

444

44



kT (keV)

- **1**

Absorbed Plasma