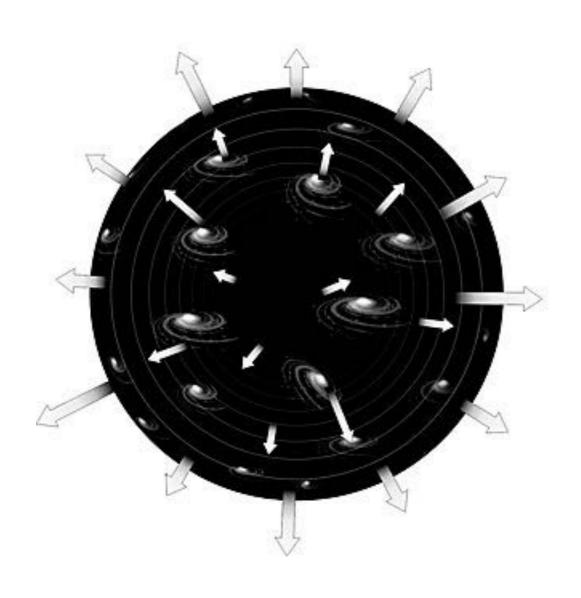
Near-infrared SN la as standard candles

Arturo Avelino

CfA, Harvard

CfA, April 24, 2018

Accelerated expansion of the Universe



Type la Supernova (SN la)



The problem

Optical samples of SN Ia for cosmology have reached their limit to constrain the nature of the dark energy (DE) because of the systematic uncertainties.

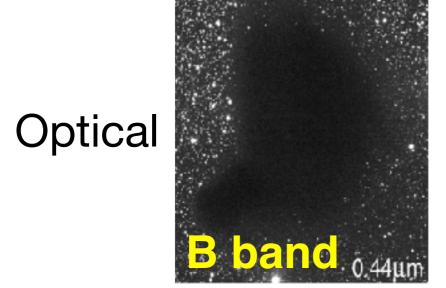
 More optical data doesn't mean better DE constraints.

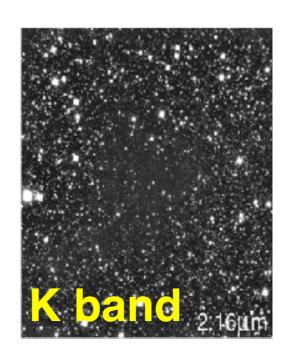
 Optical light is dimmed and reddened by dust in the host galaxy, the Milky Way, and the extragalactic medium.



A solution: NIR observations!

 Near infrared (NIR) light is much less sensitive to dust than the optical wavelengths.





Near infrared

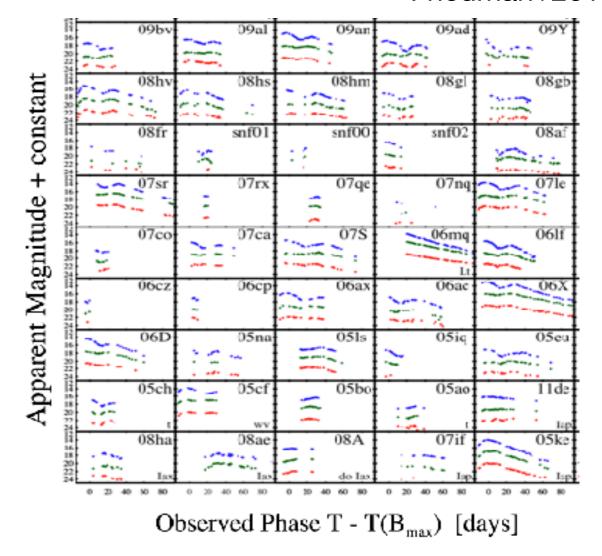
Low-z NIR sample

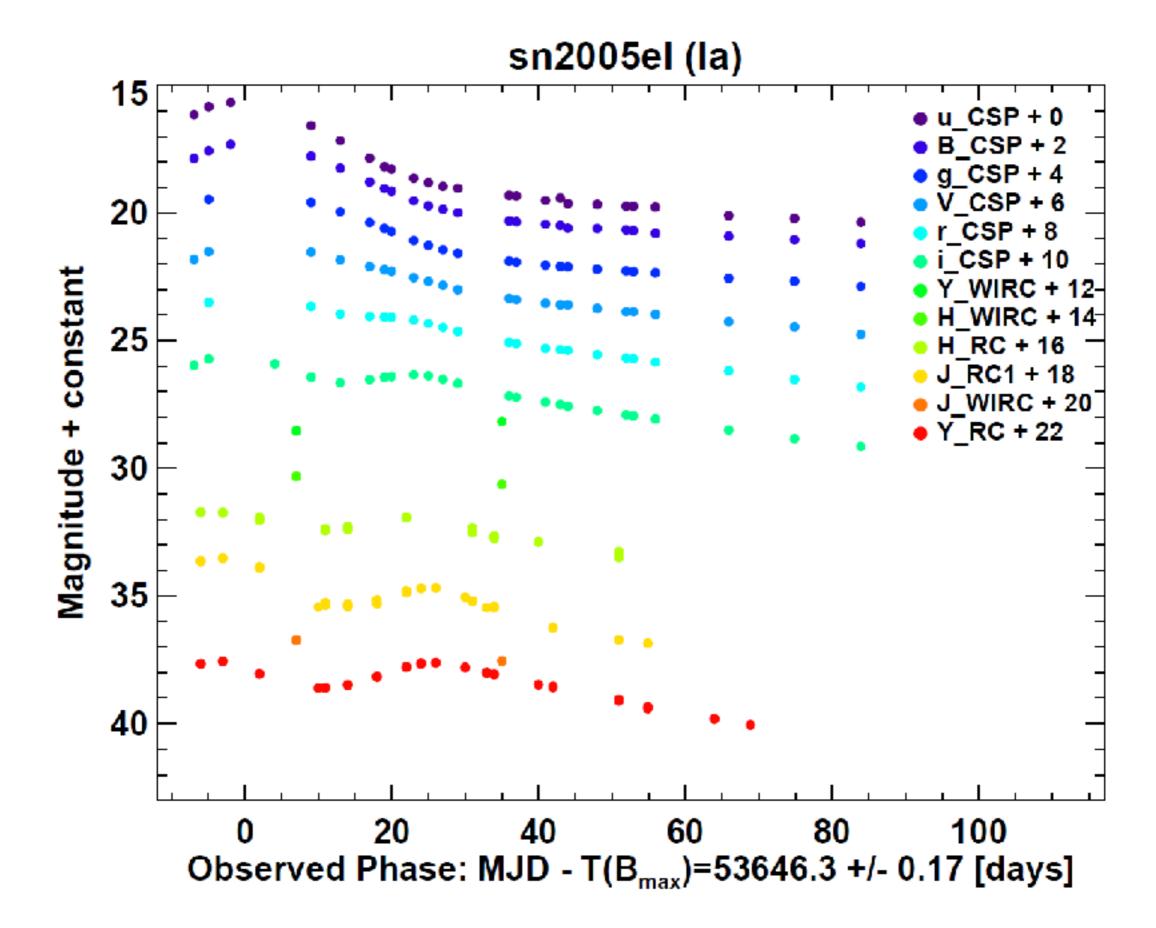
Friedman+2015

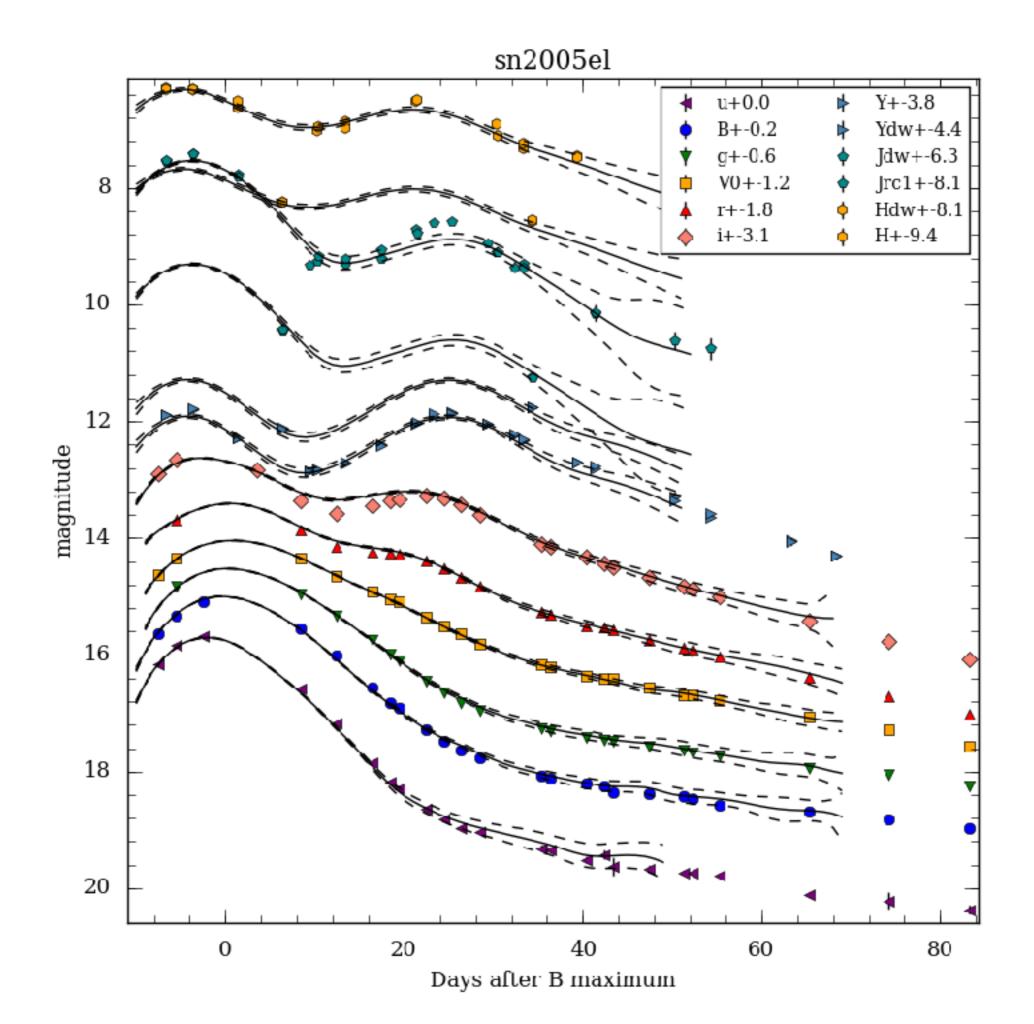
Compiled by **Andrew Friedman** (UCSD):

• CfA, CSP, Literature

 190 SN Ia with optical + NIR (YJHK) light curves







Goal

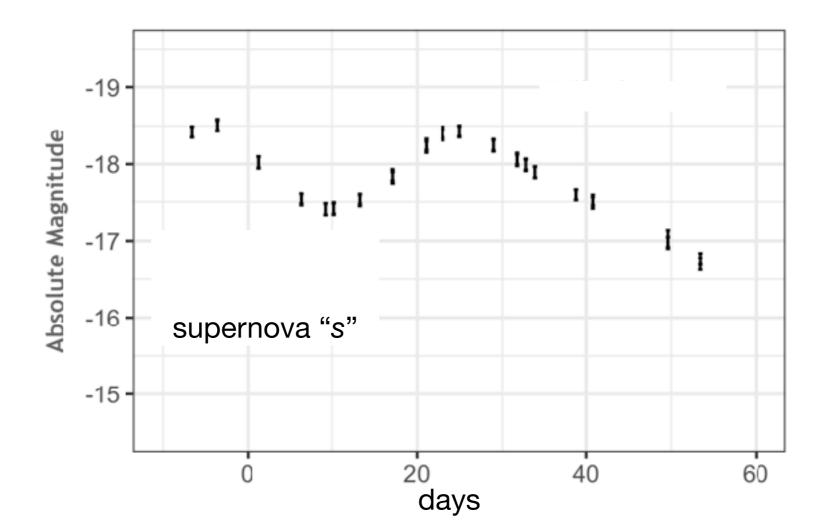
Infer the distance modulus (luminosity distance) of each SNIa from their near-infrared time-series data (aka, light curves)

Method

- ★ Construct NIR light-curve templates
 - Gaussian-Processes regression
 - Hierarchical Bayesian model
- ★ Fit the NIR light-curve template to the time series data

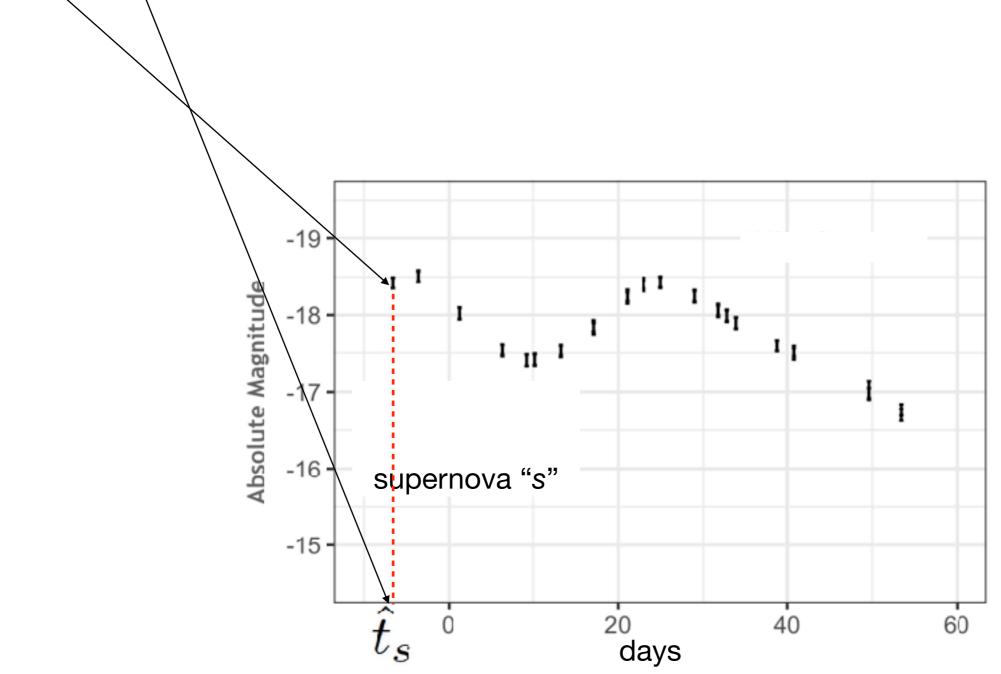
Interpolating the time series using Gaussian Processes regression

 $M_{\mathbf{r},s}(\hat{t}_s)$ = datum at a given time, for a given supernova "s" and band.



Interpolating the time series using Gaussian Processes regression

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Interpolating the time series using Gaussian Processes regression

 $M_{{f r},s}(\hat{t}_s)$ = datum at a given time, for a given supernova "s" and band.

days

$$M_{
m r,s}({f t}_s^*)|M_{
m r,s}(\hat{f t}_s),\hat{f t}_s,{f t}_s^*\sim N[ar{M}_{
m r,s}({f t}_s^*),{
m cov}(M_{
m r,s}({f t}_s^*))]$$

Interpolating the time series using Gaussian Processes regression

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days

Interpolating the time series using Gaussian Processes regression

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Mean function is computed as:

$$\mathbf{\bar{M}_{r,s}(\mathbf{t}_s^*)} = K(\mathbf{t}_s^*, \hat{\mathbf{t}}_s) \cdot \left[K(\hat{\mathbf{t}}_s, \hat{\mathbf{t}}_s) + W(\hat{\mathbf{t}}_s, \hat{\mathbf{t}}_s) \right]^{-1} \cdot \mathbf{M}_{r,s}(\hat{\mathbf{t}}_s)$$

Covariance matrix is computed as:

$$cov(\mathbf{M}_{r,s}(\mathbf{t}_s^*)) = K(\mathbf{t}_s^*, \mathbf{t}_s^*) - K(\mathbf{t}_s^*, \hat{\mathbf{t}}_s) \cdot \left[K(\hat{\mathbf{t}}_s, \hat{\mathbf{t}}_s) + W(\hat{\mathbf{t}}_s, \hat{\mathbf{t}}_s) \right]^{-1} \cdot K(\hat{\mathbf{t}}_s, \mathbf{t}_s^*)$$

Kernel:
$$K(t,t') = \sigma_K^2 \exp\left[-\frac{(t-t')^2}{2l^2}\right]$$

Notation:
$$W(\hat{t},\hat{t}')=\sigma_M^2\delta_{tt'}$$
 $\mathbf{M}_{\mathrm{r},s}(\mathbf{t}_s^*)=\{M_{\mathrm{r},s}(t_s^*)\}$ $\mathbf{M}_{\mathrm{r},s}(\hat{\mathbf{t}}_s)=\{M_{\mathrm{r},s}(\hat{t}_s)\}$

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Kernel:
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Hyperparameters computed as:

$$\ln p(\{\mathbf{M}_{r,s}\}|\{\hat{\mathbf{t}}_{s}\}, \sigma_{K}, l) =$$

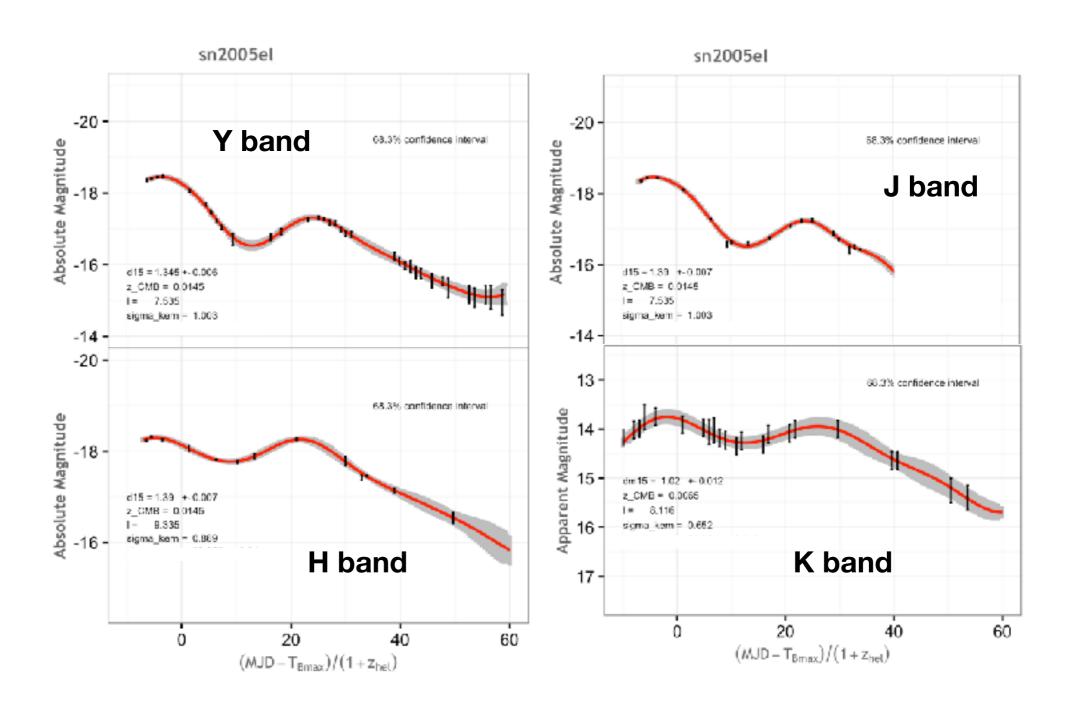
$$-\frac{1}{2} \sum_{s=1}^{N_{SN}} \{\mathbf{M}_{r,s}^{\mathsf{T}}(\hat{\mathbf{t}}_{s}) \cdot \left[K_{s}(\hat{\mathbf{t}}_{s}, \hat{\mathbf{t}}_{s}) + W_{s}(\hat{\mathbf{t}}_{s}, \hat{\mathbf{t}}_{s}) + \right.$$

$$\sigma_{\text{pec},s}^{2} I_{s} \cdot I_{s}^{\mathsf{T}} \right]^{-1} \cdot \mathbf{M}_{r,s}(\hat{\mathbf{t}}_{s}) +$$

$$\ln \left(\det \left[K_{s}(\hat{\mathbf{t}}_{s}, \hat{\mathbf{t}}_{s}) + W_{s}(\hat{\mathbf{t}}_{s}, \hat{\mathbf{t}}_{s}) + \sigma_{\mu_{\text{pec}},s}^{2} I_{s} \cdot I_{s}^{\mathsf{T}} \right] \right) +$$

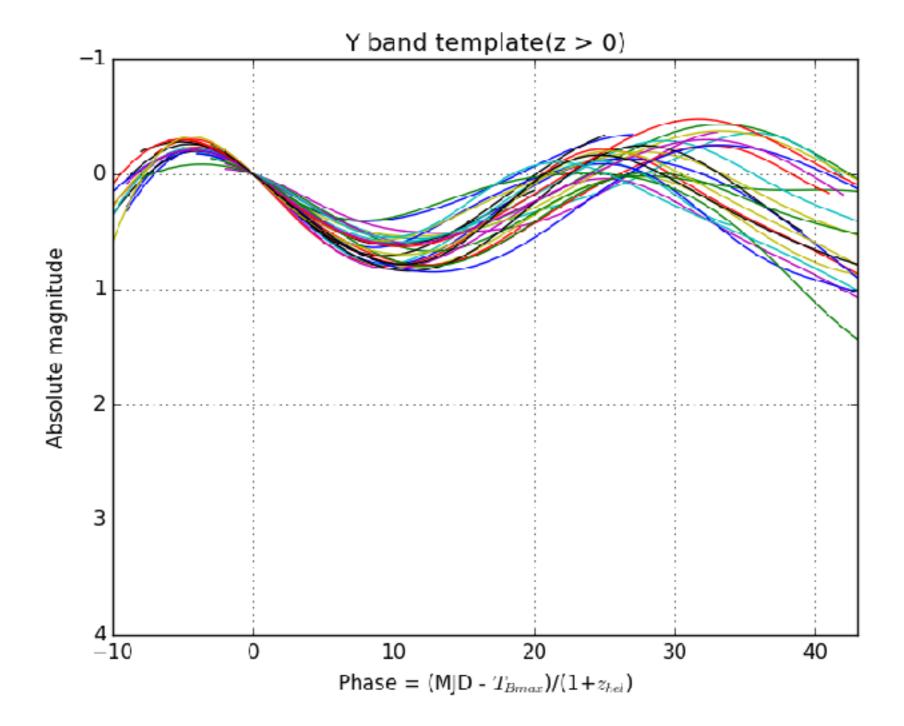
$$N_{LC,s} \ln 2\pi \}, \quad (A6)$$

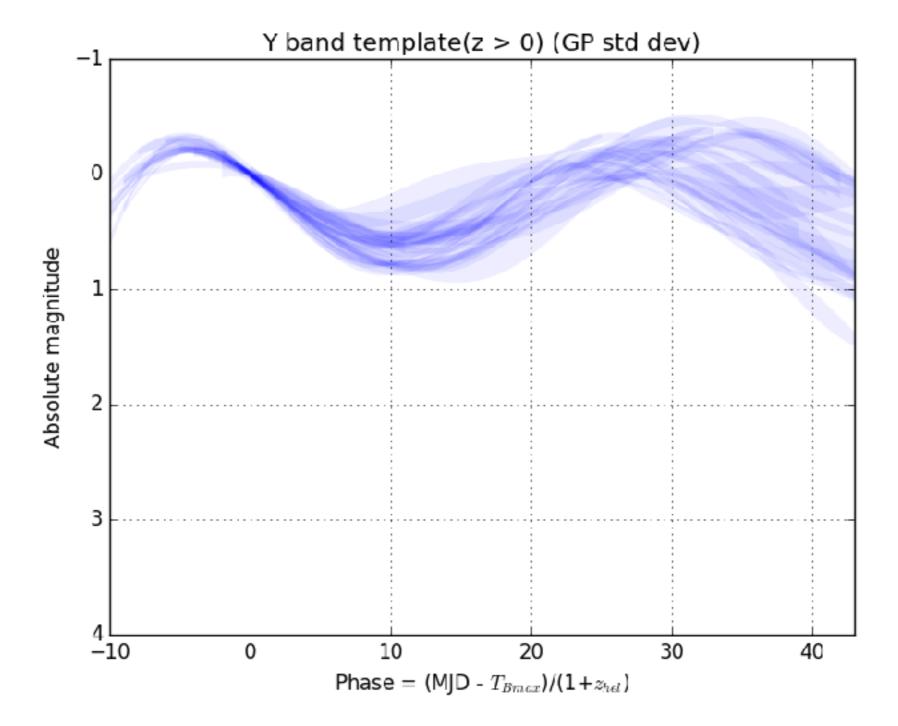
Gaussian-Process fit

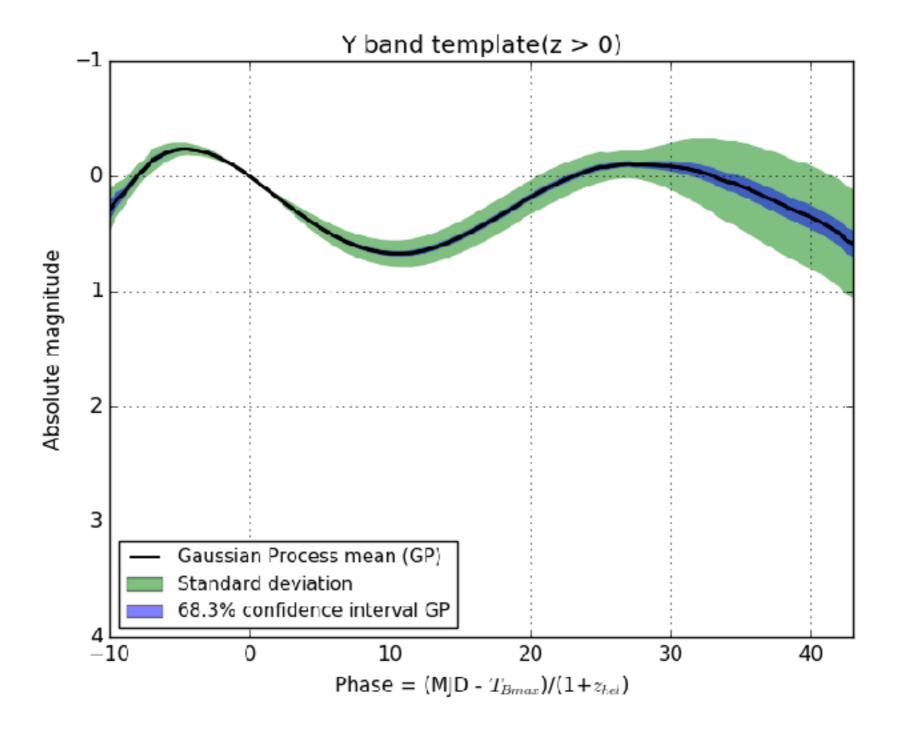


Templates

Hierarchical Bayesian model







Bayesian Hierarchical model

Constructing the NIR light-curve templates

1st level of the hierarchy:

$$\bar{M}_s \sim N(\tilde{M}_s, \sigma^2_{\bar{M},s})$$

Bayesian Hierarchical model

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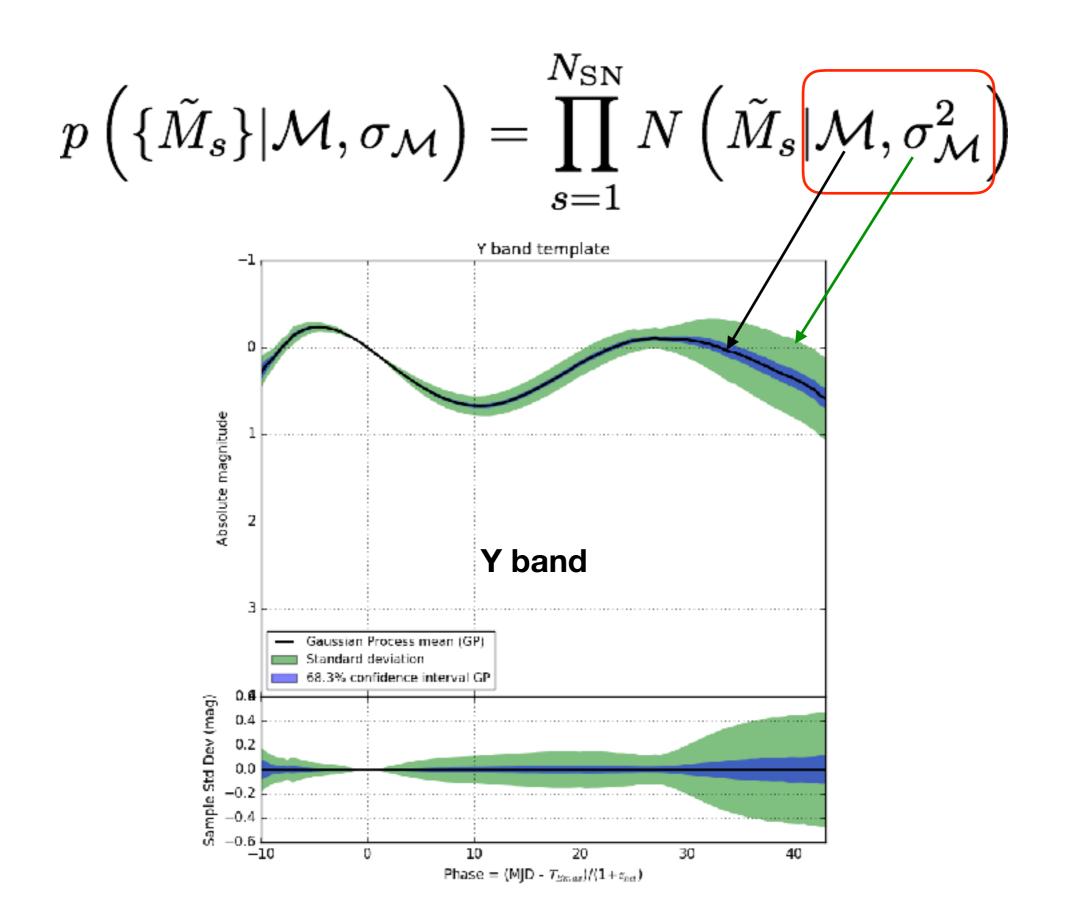
$$\bar{M}_s \sim N(\tilde{M}_s, \sigma^2_{\bar{M},s})$$

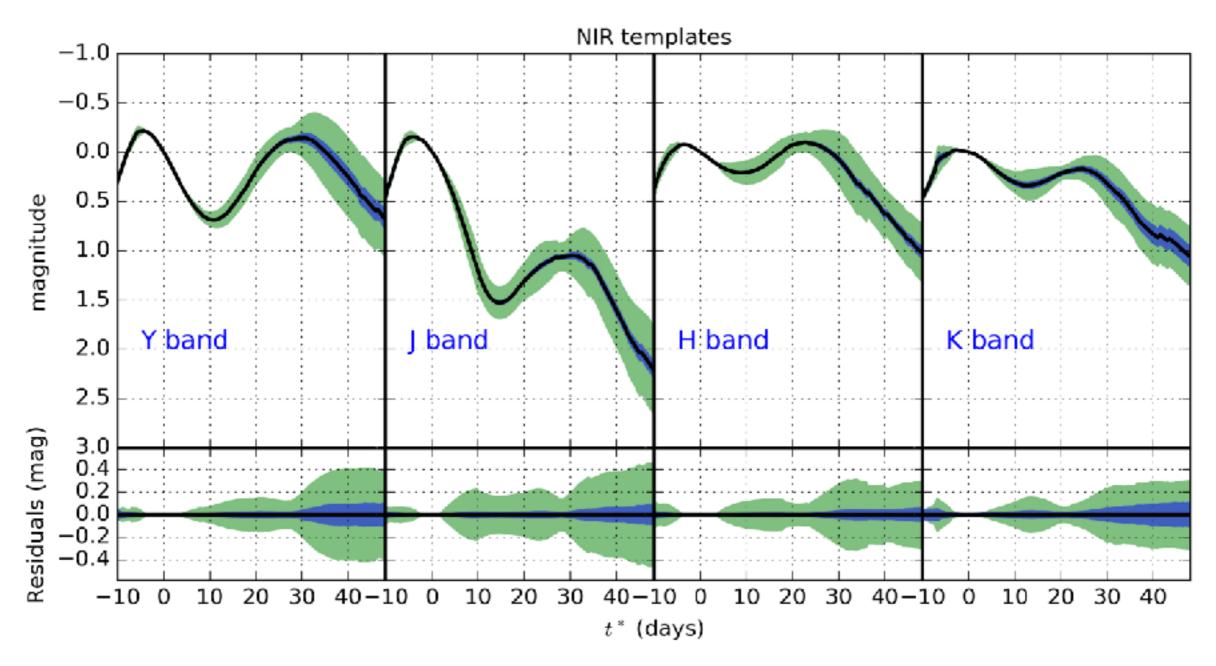
2nd level of the hierarchy:

We assume that the $\tilde{M}_s(t_*)$ are drawn from a Gaussian distribution with mean $\mathcal{M}(t_*)$ and variance $\sigma^2_{\mathcal{M}}$:

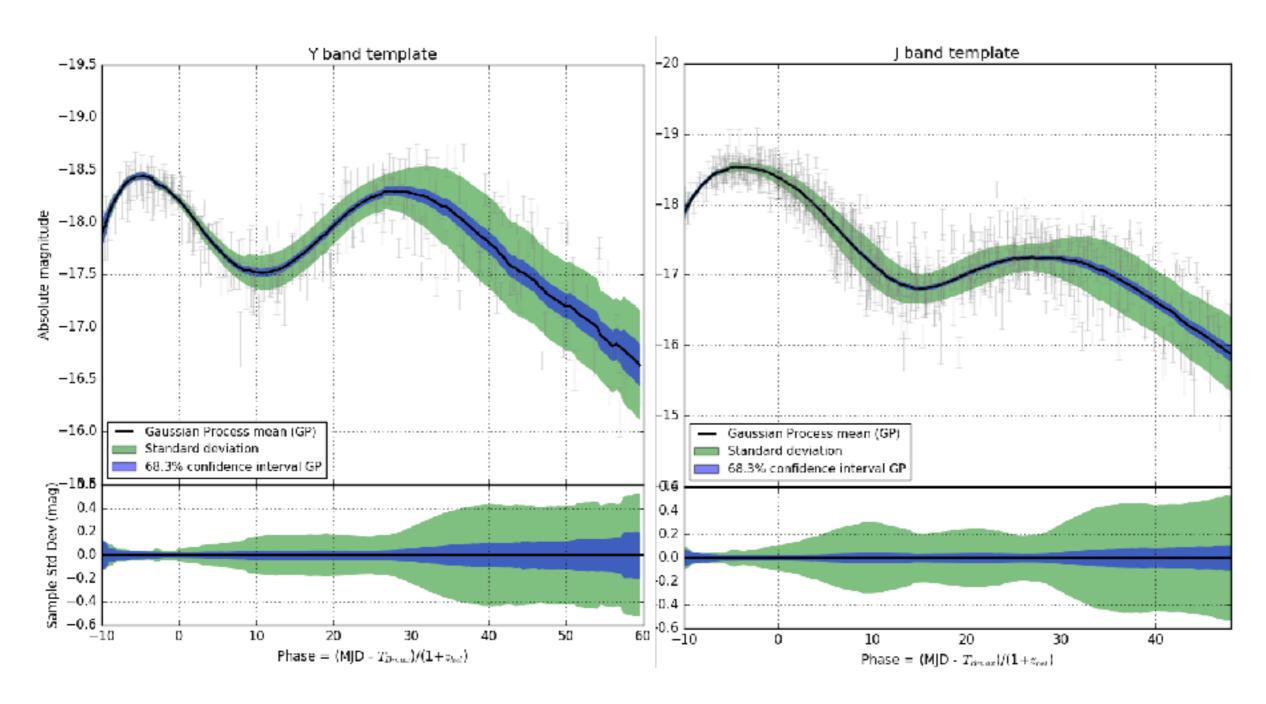
$$p\left(\{ ilde{M}_s\}|\mathcal{M},\sigma_{\mathcal{M}}
ight) = \prod_{s=1}^{N_{\mathrm{SN}}} N\left(ilde{M}_s \mathcal{M},\sigma_{\mathcal{M}}^2
ight)$$
 Template

Hierarchical Bayesian model

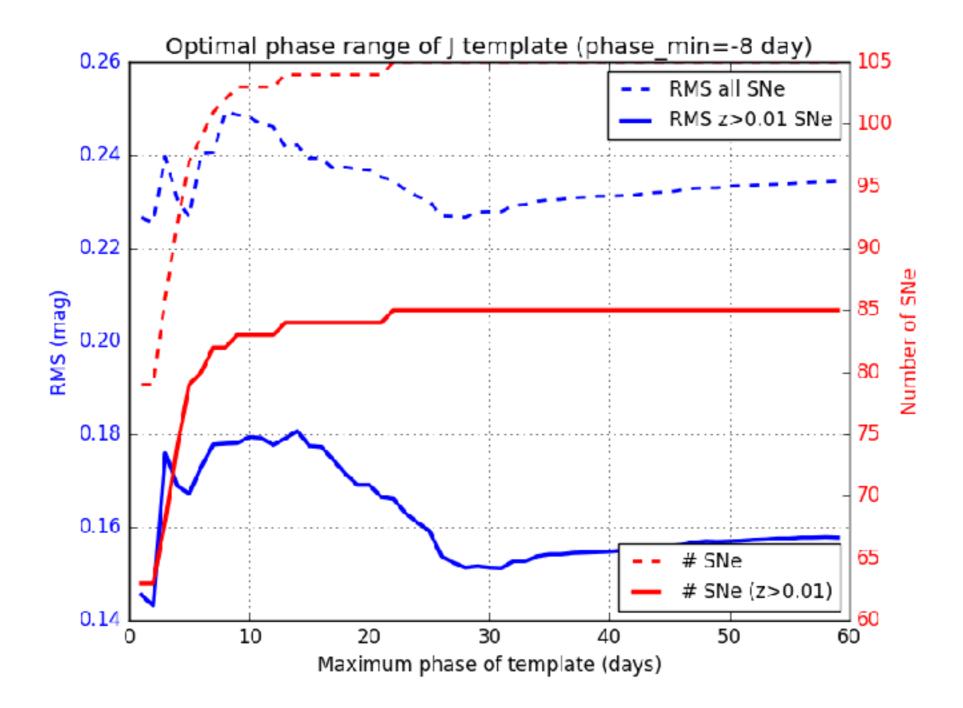


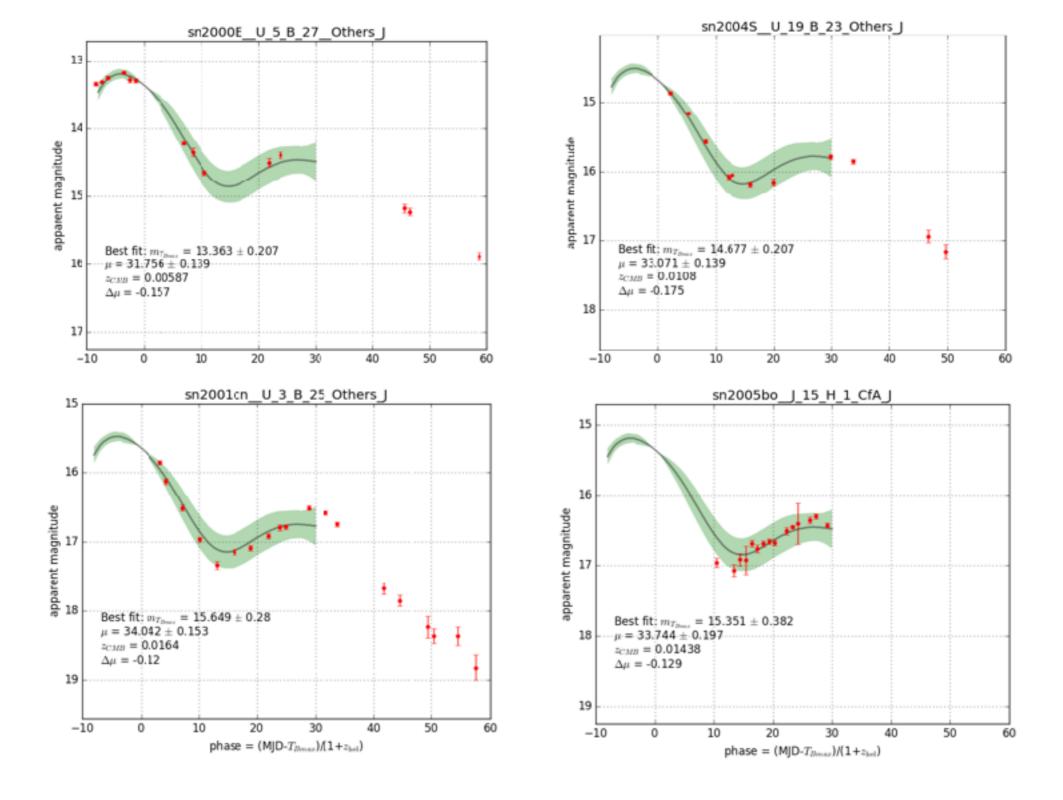


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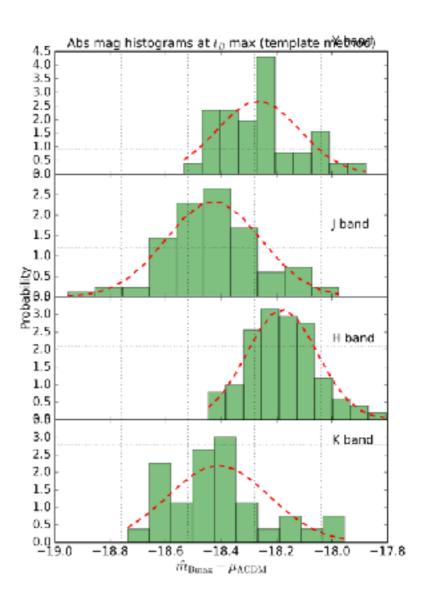
$$\Delta m_s(\hat{t}) \equiv m_s(\hat{t}) - \mathcal{M}(\hat{t}) - m_{0,s} \tag{11}$$

where $m_s(\hat{t})$ and $\mathcal{M}(\hat{t})$ are the apparent magnitude and the magnitude of the normalized template at phase \hat{t} , respectively. We can express this difference for all the $N_{\text{LC},s}$ phases in a given LC as the vector,

$$\Delta m_s \equiv \begin{pmatrix} \Delta m_s(\hat{t}_1) \\ \Delta m_s(\hat{t}_2) \\ \vdots \\ \Delta m_s(\hat{t}_{N_{\rm LC},s}) \end{pmatrix}. \tag{12}$$

Then, to determine $m_{0,s}$ we minimize the negative of the log likelihood function $L(m_{0,s})$ defined as

$$-2\ln L(m_{0,s}) = \Delta \boldsymbol{m}_s^{\top} \cdot C^{-1} \cdot \Delta \boldsymbol{m}_s$$
 (13)



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$$C_{ij} \equiv \text{Cov} \left(\Delta m_s(\hat{t}_i), \Delta m_s(\hat{t}_j) \right)$$
 (14)

$$= \sigma_{\mathcal{M}}(\hat{t}_i) \sigma_{\mathcal{M}}(\hat{t}_j) \exp \left[-\frac{(\hat{t}_i - \hat{t}_j)^2}{2l^2} \right] +$$

$$\hat{\sigma}_{m,s}^2(\hat{t}_i) \delta_{ij}$$
 (15)

where $\sigma_{\mathcal{M}}(\hat{t})$ is the population standard deviation of the sample distribution of magnitudes at time \hat{t} , determined from Eq. (B2) during the training process used to construct the mean LC template, with the hyperparameter l computed via Eq. (A6), while $\hat{\sigma}_{m,s}^2(\hat{t}_i)$ is the photometric error of the datum $m_s(\hat{t}_i)$.

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From Eq. (13), we can calculate an analytic expression for the maximum likelihood estimator (MLE) of the apparent magnitude at B-band maximum light, $\hat{m}_{0,s}$, given by:

$$\hat{m}_{0,s} = \left[\sum_{i,j}^{N_{LC,s}} (C^{-1})_{ij} \right]^{-1} \times \sum_{i}^{N_{LC,s}} \left[\left(m_s(\hat{t}_i) - \mathcal{M}(\hat{t}_i) \right) \sum_{j}^{N_{LC,s}} (C^{-1})_{ij} \right], \quad (16)$$

with the MLE of the uncertainty of $\hat{m}_{0,s}$ given as

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$$\mu_s = \hat{m}_{0,s} - \langle M_0 \rangle \tag{19}$$

with uncertainty given as

$$\sigma_{\mu,s} = \sqrt{\sigma_{0,s}^2 + \sigma_{\text{int}}^2} \tag{20}$$

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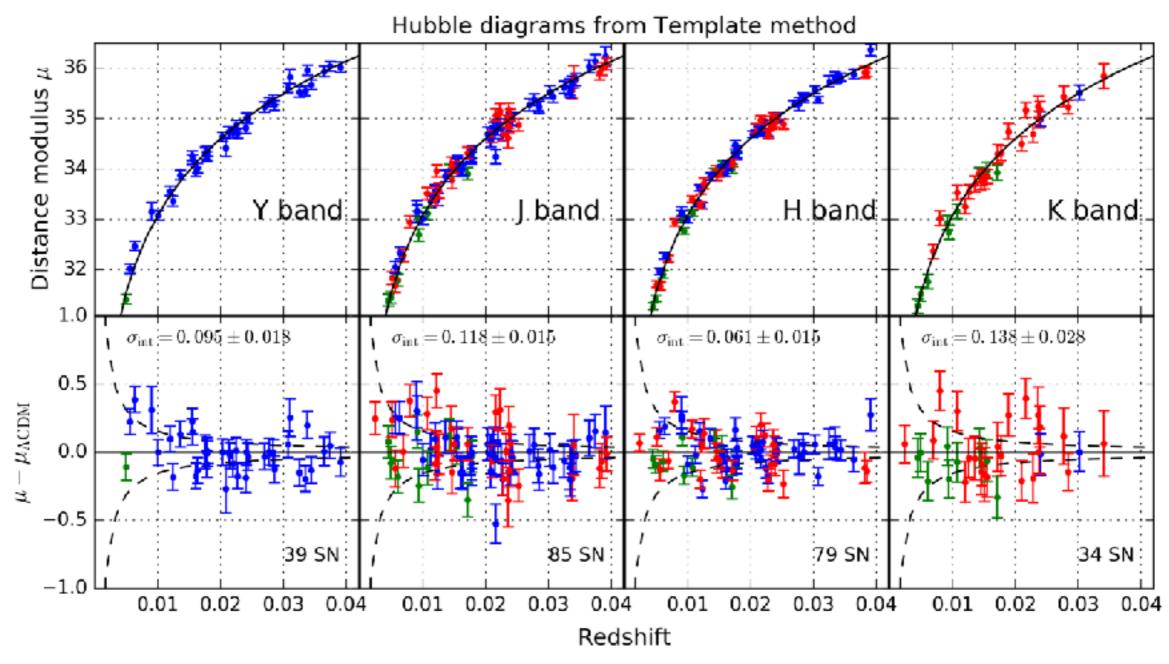
Intrinsic dispersion

Scatter in the Hubble residuals after accounting for peculiar-velocity and photometric uncertainties.

Intrinsic dispersion $\sigma_{\rm int}$:

$$-2\ln\mathcal{L}(\sigma_{\rm int}^2) = \sum_{s}^{N_{\rm SN}} \left[\ln\left(\sigma_{0,s}^2 + \sigma_{\rm int}^2 + \sigma_{\mu_{\rm pec},s}^2\right) + \frac{\delta\mu_s^2}{\sigma_{0,s}^2 + \sigma_{\rm int}^2 + \sigma_{\mu_{\rm pec},s}^2} \right]$$

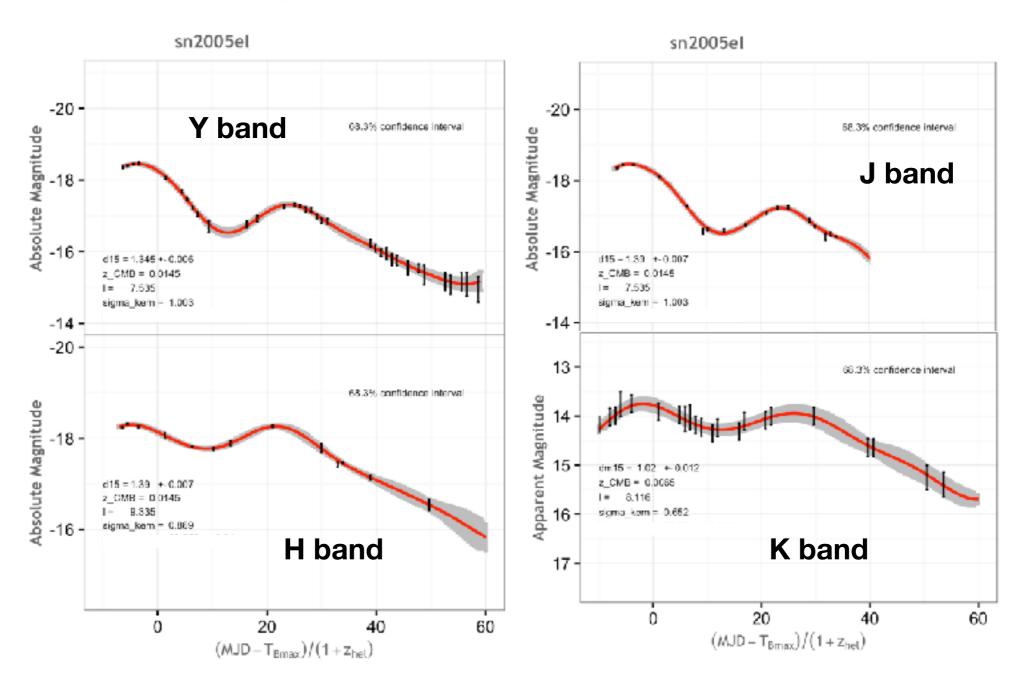
Blondin, Mandel, Kirshner, 2011



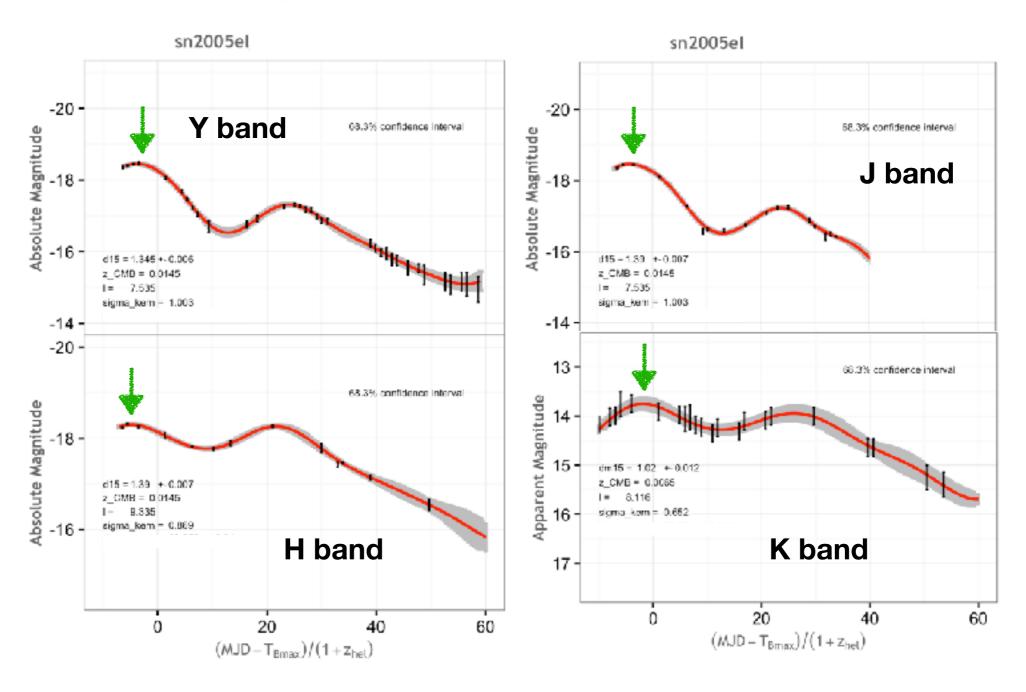
Arturo Avelino, "Near-infrared SN la as standard candles"

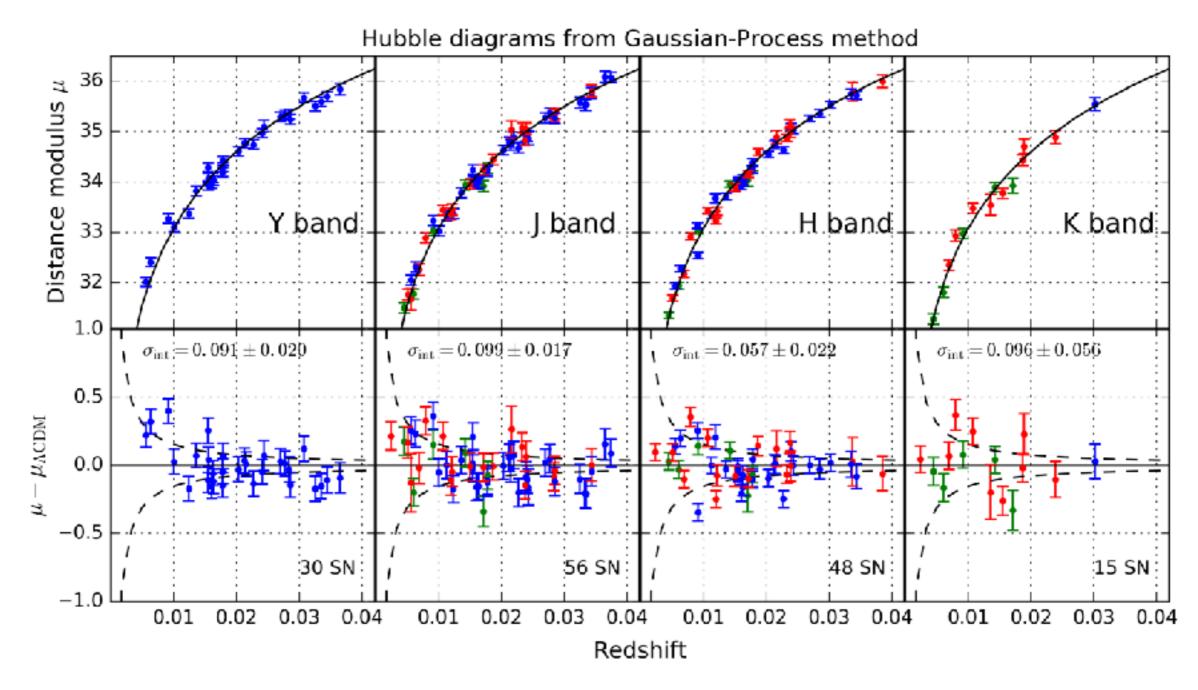
Gaussian-Process method

Gaussian-Process Method



Gaussian-Process Method





Arturo Avelino, "Near-infrared SN la as standard candles"

Combining multiple NIR bands

Distance modulus

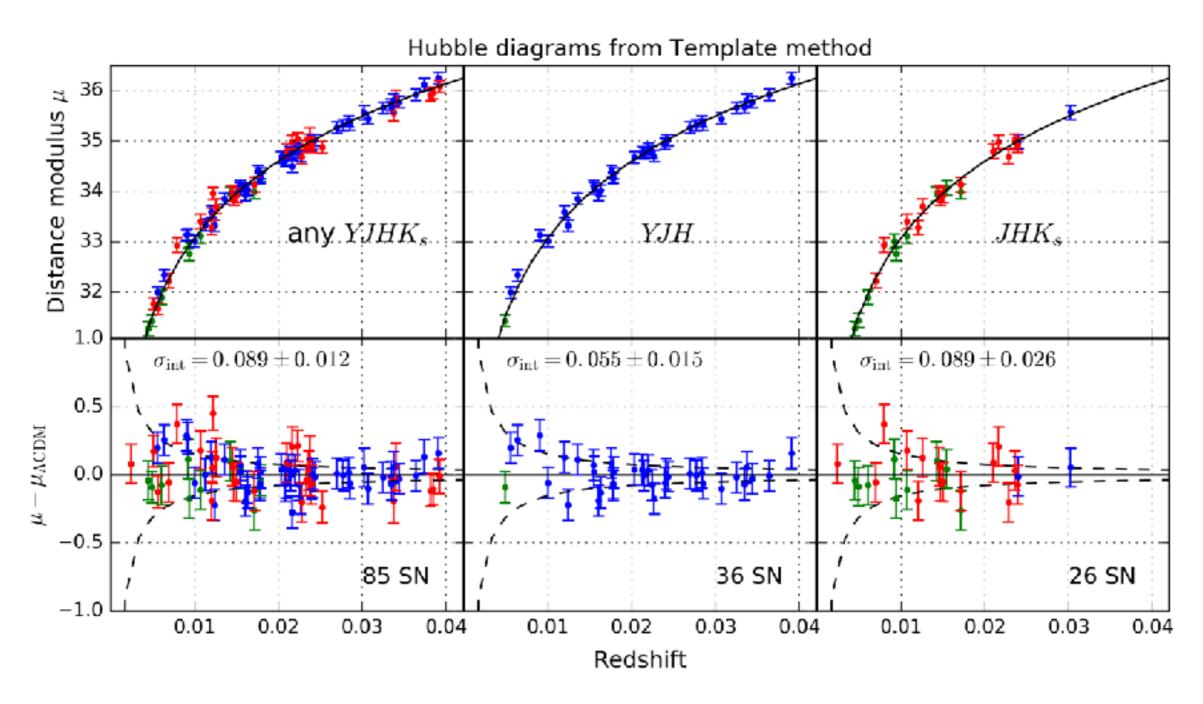
4.3. Distance modulus from the combined NIR bands

From the distance moduli $(\mu_s^Y, \mu_s^J, \mu_s^H, \mu_s^K)$ for a given supernova s determined from each NIR band following either of the two methods described above, we determine the "total" distance modulus $\hat{\mu}_s$ in each method. First we define the vector of residuals

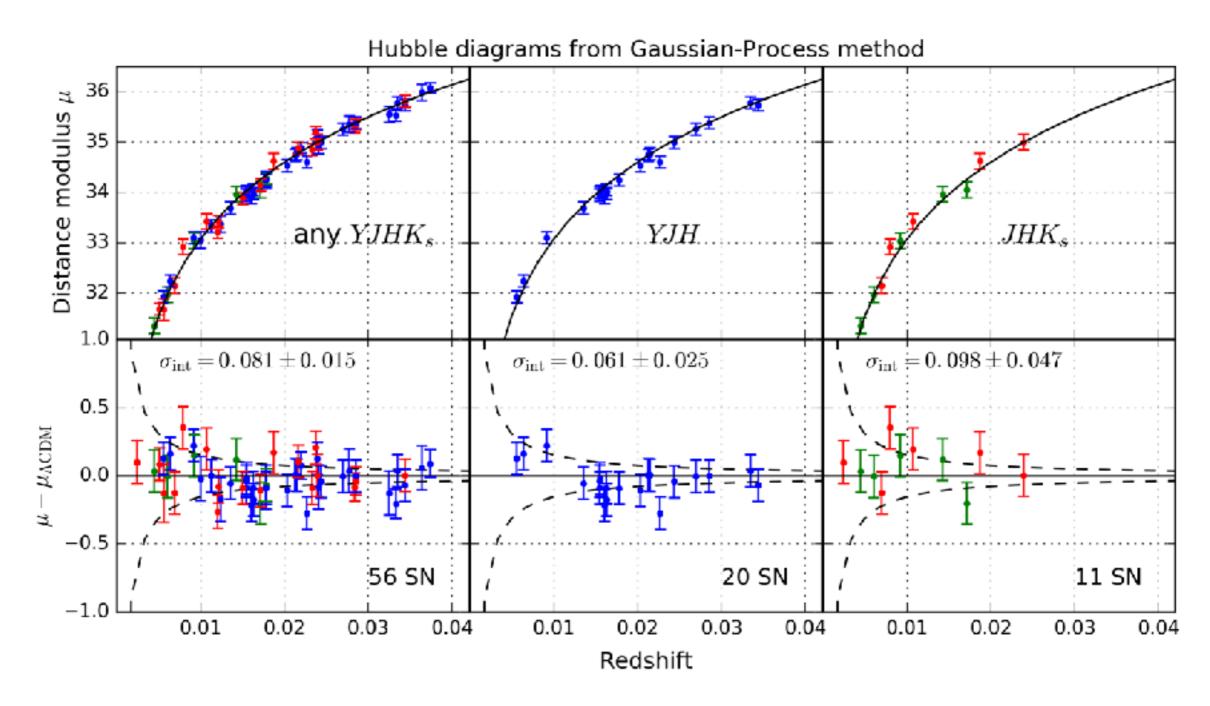
$$\delta \boldsymbol{\mu}_{s} \equiv \begin{pmatrix} \mu_{s}^{Y} - \hat{\mu}_{s} \\ \mu_{s}^{J} - \hat{\mu}_{s} \\ \mu_{s}^{H} - \hat{\mu}_{s} \\ \mu_{s}^{K} - \hat{\mu}_{s} \end{pmatrix}. \tag{25}$$

where μ_s^Z is given by either Eq. (19) or (23). Then, to determine $\hat{\mu}_s$ we minimize the negative of the likelihood function $L(\hat{\mu}_s)$ defined as

$$-2\ln L(\hat{\mu}_s) = \delta \boldsymbol{\mu}_s^{\top} \cdot C_{\boldsymbol{\mu}}^{-1} \cdot \delta \boldsymbol{\mu}_s \tag{26}$$



Arturo Avelino, "Near-infrared SN la as standard candles"

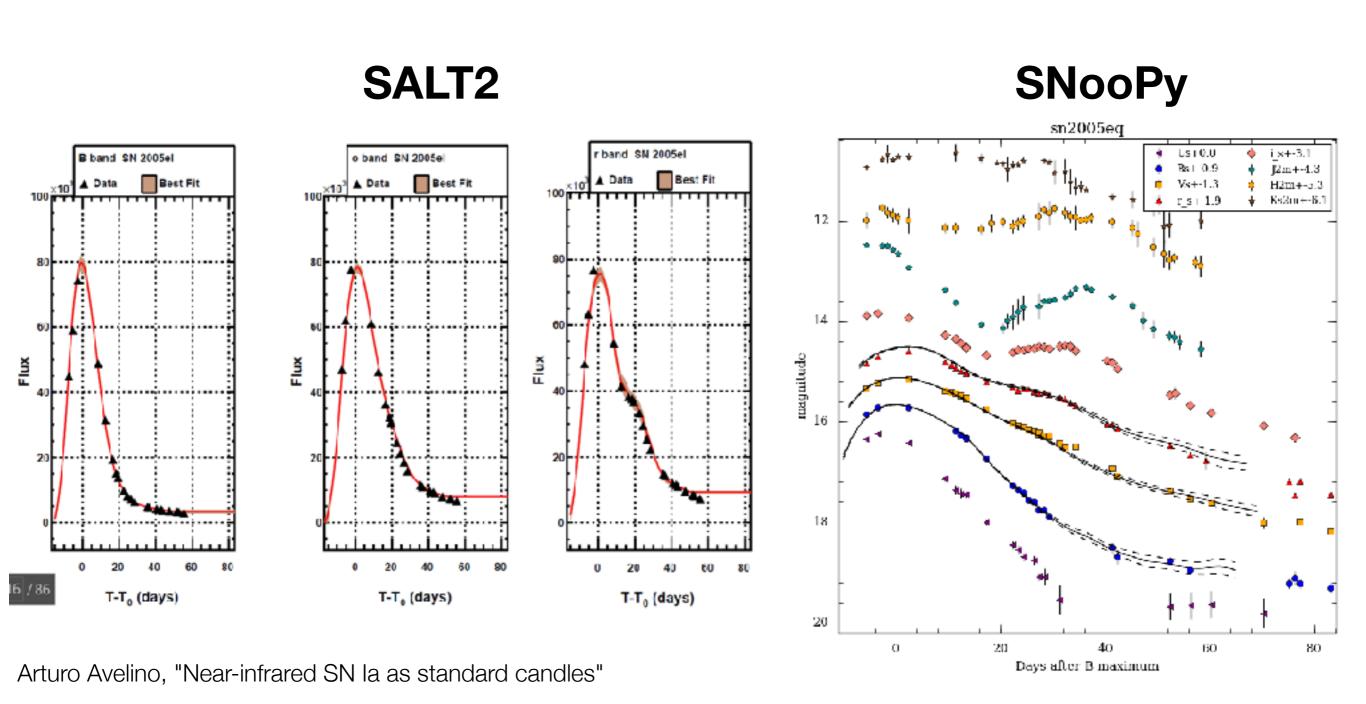


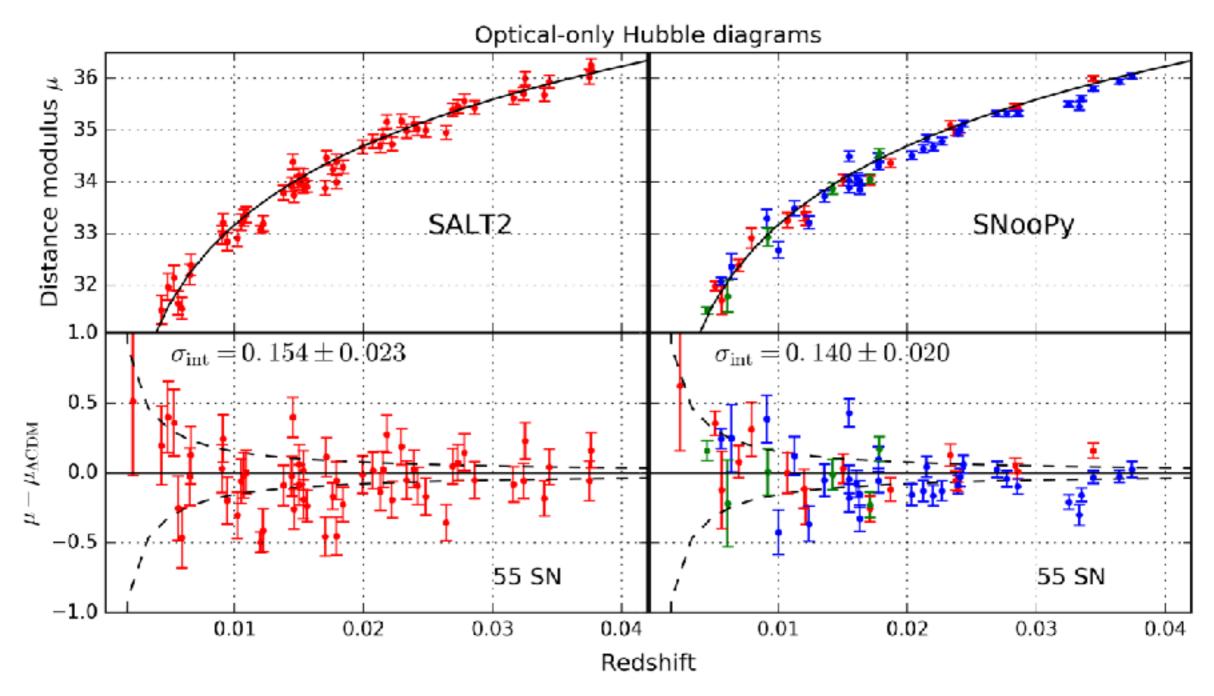
Arturo Avelino, "Near-infrared SN la as standard candles"

How good or bad are these results?

Optical Hubble diagram

Fitting the *optical* light curves only





Arturo Avelino, "Near-infrared SN la as standard candles"

Intrinsic dispersion and wRMS summary

-				
	Band	Method	$\sigma_{ m int}$	wRMS (mag)
-	Y	Template	0.095 ± 0.018	0.129
	Y	GP	0.091 ± 0.020	0.125
	J	Template	0.118 ± 0.015	0.156
	J	GP	0.099 ± 0.017	0.137
	H	Template	0.061 ± 0.015	0.113
	H	GP	0.057 ± 0.022	0.117
	K_s	Template	0.138 ± 0.028	0.180
	K_s	GP	0.096 ± 0.056	0.170
	any $YJHK_s$	Template	0.089 ± 0.012	0.123
	any $YJHK_s$	GP	0.081 ± 0.015	0.118
	YJH	Template	0.055 ± 0.015	0.097
	YJH	GP	0.061 ± 0.025	0.105
	JHK_s	Template	0.089 ± 0.026	0.134
	JHK_s	GP	0.098 ± 0.047	0.149
	Optical	SALT2	0.154 ± 0.023	0.216
	Optical	SNooPy	0.140 ± 0.020	0.146

RAISIN = SN IA in the IR

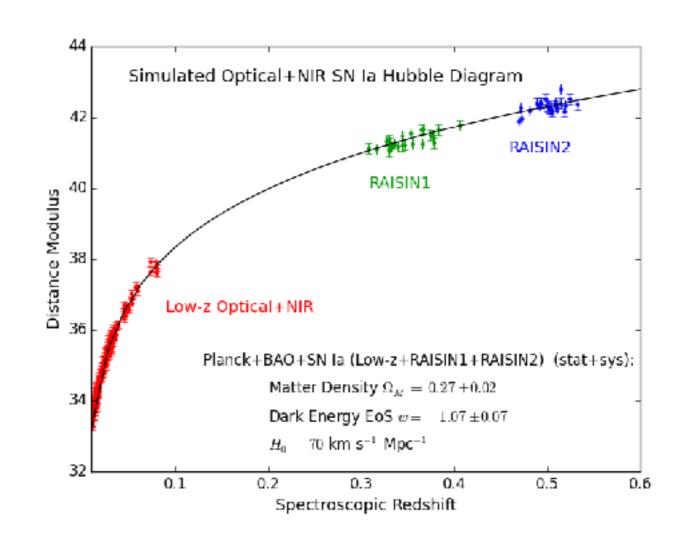
Tracing cosmic expansion with SN Ia in the Near Infrared

RAISIN-1

• 23 SN Ia, redshift ~ 0.3

RAISIN-2

• 24 SN Ia, redshift ~ 0.5



Arturo Avelino, "Near-infrared SN la as standard candles"

Take away

 NIR SN Ia are very good standard candles compared with optical observations.



Optical



 Very promising for cosmology when combining optical+NIR observations: RAISIN program, WFIRST.

